



# Measurements of tritium recycling and isotope exchange in TFTR

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## Abstract

Tritium Balmer-alpha ( $T_\alpha$ ) emission, along with  $H_\alpha$  and  $D_\alpha$ , is observed in the current D–T experimental campaign in TFTR. The data are a measure of the fueling of the plasma by tritium accumulated in the TFTR limiter and the spectral profile maps neutral hydrogenic velocities.  $T_\alpha$  is relatively slow to appear in tritium neutral beam heated discharges, ( $T_\alpha/(H_\alpha + D_\alpha + T_\alpha) = 11\%$  after 8 tritium-only neutral beam discharges). In contrast, the  $T_\alpha$  fraction in a sequence of six discharges fueled with tritium puffs, increased to 44%. Larger transient increases (up to 75%  $T_\alpha$ ) were observed during subsequent tritium gas puffs. Analysis of the Doppler broadened spectral profiles revealed overall agreement with the dissociation, charge exchange, sputtering and reflection velocities predicted by the neutral Monte-Carlo code DEGAS with some deficiency in the treatment of dissociation products in the 10–100 eV range.

*Keywords:* TFTR; Tritium inventory and economy; Line emission diagnostic; Monte Carlo simulation; Wall particle retention

## 1. Introduction

Hydrogenic ions, recycled from the limiter, are an important factor in the composition and reactivity of plasmas in the Tokamak Fusion Test Reactor, TFTR [1]. Deuterium–tritium operation of TFTR has provided a special opportunity to study tritium recycling. The recycling can be observed in Balmer-alpha emission from neutral hydrogen isotopes in the inboard plasma edge region where the plasma contacts a toroidal carbon limiter. Neutral tritium has a Balmer-alpha transition,  $T_\alpha$ , that is analogous to the  $H_\alpha$  and  $D_\alpha$  transitions long used in edge plasma diagnostics. The wavelengths of the  $H_\alpha$ ,  $D_\alpha$  and  $T_\alpha$  emission lines are separated by small isotope shifts arising from differences in the reduced mass and are at 656.28 nm, 656.104 nm and 656.045 nm, respectively. Note that the wavelength separation between  $D_\alpha$  and  $T_\alpha$  is one third of that between  $D_\alpha$  and  $H_\alpha$  and the lines are partially blended.

The first spectroscopic measurements of  $T_\alpha$  emission from a fusion plasma were made on TFTR using a Fabry–Perot interferometer [2]. The relative fraction of the hydro-

genic isotopes in the  $H_\alpha$ ,  $D_\alpha$  and  $T_\alpha$  spectral profile reflects the hydrogenic inventory in the surface of the limiter. This, in turn, is a function of the fueling history of prior discharges, burial by codeposition and the mobility of hydrogen isotopes in graphite. Spectral resolution of the line profile gives information on processes in the edge plasma. The emission is Doppler shifted and the spectral line profile maps the velocity distribution of the neutral hydrogen isotopes. Since the edge density is insufficient to thermalize the velocity distribution, the contributions of the various reaction pathways that generate hydrogenic atoms may be identified with the different wavelength regions (velocities) in the spectral profile. Previous studies of neutral velocities have been made on TEXTOR [3,4] and DITE [5]. In this paper we will describe spectrally resolved measurements of the  $H_\alpha$ ,  $D_\alpha$  and  $T_\alpha$  emission on TFTR and compare the velocities to those predicted by the neutrals code DEGAS. We also compare the isotope changeover rate in the Balmer-alpha emission in a sequence of discharges with tritium neutral beam injection to a sequence with tritium gas puffs. Modeling of tritium recycling in discharges with deuterium-only neutral beam injection is reported in Ref. [6] and an account of tritium retention, lithium conditioning and advanced tokamak regimes on TFTR is given in Ref. [7] in these proceedings.

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## 2. Experimental arrangement

TFTR plasmas have a circular cross section with a plasma boundary defined by an inboard toroidal belt limiter composed of carbon composite tiles in high heat flux regions as well as graphite tiles, both supported by water cooled Inconel 718 backing plates. The limiter experiences erosion, codeposition of hydrogen with carbon and neutron flux from the DT plasmas. After exposure to many plasma discharges, each limiter sector develops a poloidally asymmetric 'footprint' or eroded area, surrounded by areas of net codeposition [8]. The rise in bulk limiter temperature remains below 50°C during a discharge [9], however infrared camera measurements indicate that the surface temperature in localized hot spots increases up to 1000°C or more depending on the discharge conditions and auxiliary heating power [10].

Light from TFTR is collected by a telescope and transmitted via a fiber optic to a remote Fabry Perot interferometer. The telescope views a region on the TFTR inboard limiter 20 cm in diameter at the midplane. The Zeeman effect splits the  $H_{\alpha}$ ,  $D_{\alpha}$  and  $T_{\alpha}$  lines into an unshifted  $\pi$  component, polarized parallel to the field direction and two  $\sigma$  components polarized perpendicular to the field. Previous measurements of the Zeeman splitting showed that the location of the emission region was close to the inner limiter. For the present measurements, a polarizing filter was placed in front of the telescope lens and oriented to transmit only the unshifted  $\pi$  component. The light was recollimated and transmitted through a prefilter to select a 7 Å band centered at the  $D_{\alpha}$  wavelength. It was then spectrally analyzed with a scanning Fabry-Perot interferometer with a wavelength resolution of 0.23 Å. Alignment of the interferometer was maintained by an electronic controller locked to helium neon laser light transmitted by the interferometer between plasma discharges [11].

## 3. Tritium recycling

The fraction of tritium in the hydrogenic recycling was surprisingly low in discharges fueled by tritium neutral beam injection. At low tritium concentrations the  $T_{\alpha}$  line appears as a displacement or 'bulge' in the short wavelength side of  $D_{\alpha}$  and the tritium fraction in the total hydrogenic emission,  $T_{\alpha}/(H_{\alpha} + D_{\alpha} + T_{\alpha})$ , is revealed by subtraction of a companion deuterium discharge observed before tritium injection. In this way,  $T_{\alpha}$  fractions down to 2% may be detected. In the first discharge with mixed tritium/deuterium neutral beam injection, (December 1993), the level of tritium recycling was too low to be detectable (< 2%). After several deuterium/tritium neutral beam fueled discharges the  $T_{\alpha}$  feature appeared and increased up to a level of 7.5% [2]. In a subsequent H-mode study [12] there were 9 sequential discharges with tritium-

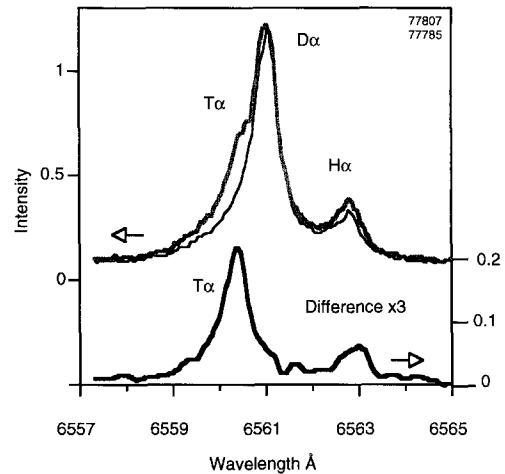


Fig. 1. Tritium-alpha emission is evident on the short wavelength wing of the Balmer-alpha spectral profile (thick line). A prior deuterium companion discharge (thin black line) is also shown and the difference, magnified by a factor 3 represents the  $T_{\alpha}$  profile. A small increase in  $H_{\alpha}$  is also evident. The  $T_{\alpha}$  feature accounts for 9% of the total area.

only neutral beam fueling intended to maximize the tritium fraction in the plasma. The injected tritium averaged 44 Ci/discharge and the  $T_{\alpha}$  fraction increased by approximately 1% each discharge. The spectral profile from the penultimate discharge in the sequence is shown in Fig. 1. In contrast, strong tritium fueling through a gas puff produced a large  $T_{\alpha}$  fraction. Fig. 2 shows such a spectral profile dominated by the  $T_{\alpha}$  peak, with resolved peaks due to  $H_{\alpha}$ ,  $D_{\alpha}$  and  $T_{\alpha}$ . The isotope changeover is discussed below.

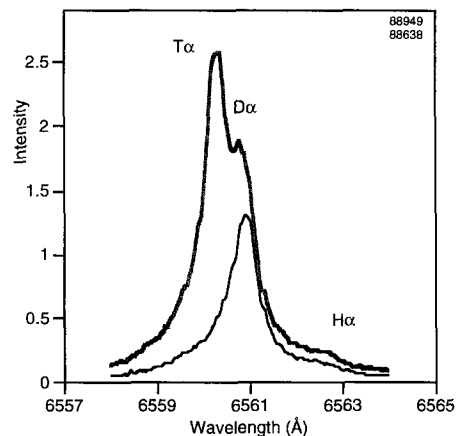


Fig. 2. Spectral profiles from discharges (i) with a tritium gas puff (thick line) and (ii) a prior D companion discharge with a deuterium puff (thin line).  $T_{\alpha}$  accounts for 54% of the total area of (i).

#### 4. Spectral analysis

The tritium capability of TFTR has provided an important opportunity to compare the velocity distribution of tritium and deuterium. To first order, one expects the tritium velocity to be lower by a factor  $\sqrt{2/3}$  due to the mass difference, however there are other possible effects. During dissociation, the lower velocity of the tritium atoms can influence the pathways of molecular breakup. For example, Higo [13] found evidence for isotope effects originating from the different widths of the Frank–Condon region, the difference in the nuclear mass and different probabilities for curve crossing. In addition, since the tritium neutrals are moving more slowly they will experience, on average, a somewhat different plasma environment to deuterium.

The velocity distribution of neutral tritium and deuterium were compared by spectral analysis. The deuterium contribution to the blended  $H_\alpha$ ,  $D_\alpha$  and  $T_\alpha$  profile in Fig. 2 was subtracted in the following way. The line profile was first fit to 6 trial Gaussians, 2 for each isotope, plus a first order background. The number of free parameters was reduced from 20 to 8 by using the known wavelength differences and the mass scaling factor. The results showed the relative area due to each isotope in the composite spectral profile. A companion discharge with deuterium puffing had the same plasma current, major radius and similar RF heating power and electron density. The deuterium spectral profile (also shown in Fig. 2) was normalized to match the deuterium area derived for the DT discharge. This profile was then subtracted, yielding a profile representing the tritium-alpha line. To reveal differences in the energy distribution between tritium and deuterium, the wavelength scale was then multiplied by the mass factor,  $\sqrt{3/2}$  and the tritium-alpha profile superposed on the deuterium profile observed from the companion discharge (Fig. 3). Differences in the velocity distribu-

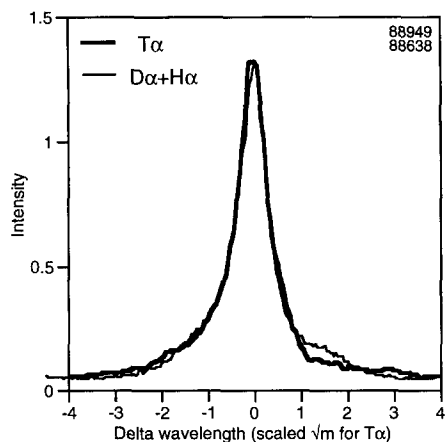


Fig. 3. Superposition of  $T_\alpha$  and  $D_\alpha + H_\alpha$  spectral profiles in normalized wavelength units.

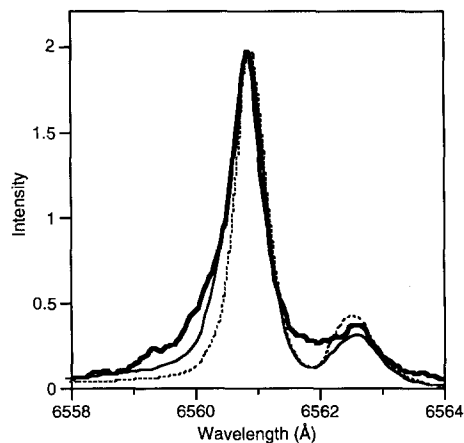


Fig. 4. Comparison of observed  $D_\alpha$  spectral profile (thick line) with the predictions of the neutrals code DEGAS (thin line). An earlier DEGAS simulation without explicit treatment of sputtering is shown (dotted line).

tion beyond the mass scaling factor should appear as differences in the shape of the spectral lines. It can be seen that the profiles are very similar. There is a trace of hydrogen present in the deuterium profile, otherwise isotopic differences beyond the  $\sqrt{3/2}$  factor are too small to resolve. This result verifies an underlying assumption in the treatment of different hydrogen isotopes in neutral codes.

The deuterium spectral profile has been used to benchmark the atomic, molecular and surface physics used in the Monte-Carlo neutral transport code DEGAS [14]. This code tracks the ‘flights’ of neutral test particles in the TFTR geometry and plasma conditions as they undergo ionization, charge exchange, dissociation and surface interactions. The plasma conditions within the last closed flux surface are input from the TRANSP [15] analysis code. The temperature and density in the scrape-off region are derived from the poloidal distribution of the Balmer-alpha emission. An extensive suite of atomic, molecular and surface physics data is used. Eight different pathways for molecular dissociation, each with differing product energies, are treated explicitly. The mass scaling factor is used in calculating the velocities of the different hydrogen isotopes, all the other atomic data are taken to be independent of isotopic species. When, through dissociation, excitation or charge exchange, a test particle emits a Balmer-alpha photon in a volume corresponding to that observed experimentally, the particle velocity is logged and, over thousands of flights, a spectral line profile built up. This predicted profile is then compared with the observed profile.

A comparison between an observed  $D_\alpha$  and  $H_\alpha$  spectral profile and the predictions of DEGAS are shown in Fig. 4. The initial disagreement in the short wavelength line wing was narrowed considerably by the explicit addition of

sputtering processes to DEGAS. Currently there is reasonable agreement between the measured and simulated spectra over a range of plasma conditions validating the treatment of charge exchange, molecular dissociation, surface reflection and sputtering in DEGAS. However a residual deficiency of neutrals in the 10–100 eV region indicates a need for further investigation of the dissociation product energies for electron temperatures in the range 100–1000 eV.

## 5. Isotope changeover

Hydrogen/deuterium isotope changeover has been previously studied on TFTR [16–18] and the number of discharges required for isotope changeover found to depend on the limiter condition. Isotope changeover has also been studied on JET [19]. The release of tritium was related to a model that incorporated implantation into a thin surface layer as well as diffusion into and out of the bulk material [20]. Early in the TFTR D–T campaign there was a period of several operation days of high power deuterium and tritium neutral beam injection followed by an extended period without tritium NBI. The rate of decay of tritium was tracked by both measurements of  $T_\alpha$  and the D–T neutrons produced by recycled tritium. The initial rate of decay fit an exponential relation with a 7.5 discharge decay constant. After an initial decrease the decay slowed markedly to a decay constant of 404 discharges [21].

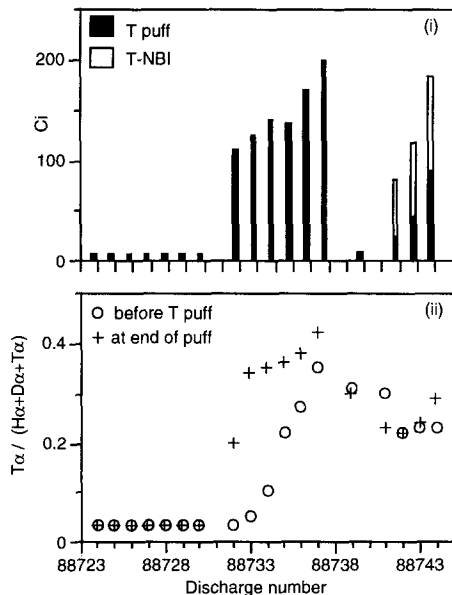


Fig. 5. (i) History of tritium injection by gas puff and neutral beams and (ii) the corresponding increase in tritium fraction in hydrogenic recycling on 1st September 1995.

Tritium gas puffs were used to generate L-mode plasmas with a maximal fraction of tritium. Fig. 5 shows the quantity of tritium injected and the corresponding rise in the tritium fraction of hydrogenic recycling on the first day of this campaign. In the first discharge with a strong tritium puff, 111 Ci were puffed into the vacuum vessel. During this puff the  $T_\alpha$  fraction increased up to 20%. This increase was temporary, after the puff the  $T_\alpha$  fraction dropped to approximately 5%. It is likely that this initial rise was due to excitation of the incoming gas stream and not due to recycled tritium. With continued tritium puffs, the  $T_\alpha$  fraction increased throughout the duration of the discharge. The maximum  $T_\alpha$  fraction observed was 75% during a subsequent 460 Ci tritium puff. However, after 8 days and a total of 10,600 Ci tritium puffed, the level of  $T_\alpha$  that persisted without a simultaneous tritium puff was approximately 50%. Although the tritium/deuterium exchange was incomplete, it did enable studies of isotope scaling in L-mode transport [22]. An account of the long term inventory of tritium injected and removed from the torus is given in Ref. [23].

In summary, tritium recycling has been studied in TFTR through observations of Balmer-alpha emission. The fraction of  $T_\alpha$  in the total hydrogenic emission was surprisingly low in discharges fueled by tritium neutral beam injection. In contrast, tritium fueling through a gas puff produced a large  $T_\alpha$  fraction. In these discharges the maximum isotopic fraction of tritium recycling that persisted without a simultaneous tritium puff was  $T_\alpha / (H_\alpha + D_\alpha + T_\alpha) \approx 50\%$ . The contributions of the various reaction pathways that generate hydrogen atoms may be identified with different wavelength regions (velocities) in the spectral profile. The observed spectral profile was compared in detail to the predictions of the DEGAS neutrals code and confirmed that the treatment of charge exchange, molecular dissociation, surface reflection and sputtering in DEGAS is broadly correct. However a residual deficiency of neutrals in the 10–100 eV region indicates a need for further investigation of the dissociation product energies for electron temperatures in the range 100–1000 eV.

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## References

- [1] K.M. McGuire, H. Adler, P. Alling et al., *Phys. Plasmas* 2 (1995) 2176–2188.
- [2] C.H. Skinner, H. Adler, R.V. Budny, J. Kamperschroer, L.C. Johnson, A.T. Ramsey and D.P. Stotler, *Nucl. Fusion* 35 (1995) 143–151.

- [3] U. Samm, P. Bogen, H. Hartig et al., *J. Nucl. Mater.* 162 (1989) 24.
- [4] D. Reiter, P. Bogen and U. Samm, *J. Nucl. Mater.* 196–198 (1992) 1059–1064.
- [5] S.J. Fielding, P.C. Johnson, M.J. Forrest, D. Guilhem and G.F. Matthews, *J. Nucl. Mater.* 162–164 (1989) 482.
- [6] R.V. Budny, M. Bitter, H. Duong, Q. Gao, D. Jassby, L.C. Johnson, S. von Goeler, A.T. Ramsey, C.H. Skinner and D.P. Stotler, presented at 12th Int. Conf. on Plasma Surface Interactions, May 1996, St.-Raphaël, France.
- [7] C.H. Skinner, E. Amarecu, G. Ascione et al., these Proceedings, p. 214.
- [8] B.E. Mills, D.A. Buchenauer, A.E. Pontau and M. Ulrickson, *J. Nucl. Mater.* 162–164 (1989) 343–349.
- [9] A. Janos, E.D. Fredrickson, K.M. McGuire, Y. Nagayama, D.K. Owens and E. Wilfred, *J. Nucl. Mater.* 196–198 (1992) 602–606.
- [10] A.T. Ramsey, C.E. Bush, H.F. Dylla, D.K. Owens, C.S. Pitcher and M.A. Ulrickson, *Nucl. Fusion* 31 (1991) 1811.
- [11] C.H. Skinner, D.P. Stotler, R.V. Budny, H. Adler and A.T. Ramsey, *Rev. Sci. Instrum.* 66 (1995) 646–648.
- [12] C.E. Bush, S.A. Sabbagh, S.J. Zweben et al., *Phys. Plasmas* 2 (1995) 2366.
- [13] M. Higo, S. Kamata and T. Ogawa, *Chem. Phys.* 73 (1982) 99.
- [14] D.P. Stotler, C.H. Skinner, R.V. Budny, H. Adler, A.T. Ramsey, D.N. Ruzic and R.B. Turkot, *Phys. Plasmas* 3 (1996) 4084.
- [15] R.V. Budny, M.G. Bell, A.C. Janos et al., *Nucl. Fusion* 35 (1995) 1497–1508.
- [16] P.H. LaMarche, H.F. Dylla, P.J. McCarthy and N.M. Ulrickson, *J. Vac. Sci. Technol.* 4 (1986) 1198.
- [17] M. Ulrickson, M.G. Bell, R. Budny et al., IAEA-CN-50/G-III-1, *Nucl. Fusion (Suppl.)* 111 (1989) 419–427.
- [18] S.J. Kilpatrick, R. Nygren, W.R. Wampler et al., *J. Vac. Sci. Technol. A* 7 (1989) 1087–1091.
- [19] L.D. Horton, P. Andrew, G. Bracco et al., *J. Nucl. Mater.* 196–198 (1992) 139–142.
- [20] P. Andrew, J.P. Coad, J. Ehrenberg et al., *Nucl. Fusion* 33 (1993) 1389.
- [21] P.C. Efthimion, L.C. Johnson, C.H. Skinner et al. and the TFTR Group, Proc. 15th Int. Atomic Energy Agency Conf. on Plasma Physics and Controlled Nuclear Fusion Research 1994, Seville, Spain, Sept. 1994, Vol. 1 (IAEA, Vienna, 1995) pp. 289–305.
- [22] S. Scott, R. Bell, C. Bush et al., Proc. 16th Int. Atomic Energy Agency Conf. on Plasma Physics and Controlled Nuclear Fusion Research, 1996, IAEA F1-CN-64/A6-6.
- [23] C.H. Skinner, W. Blanchard, J. Kamperschroer et al., *J. Vac. Sci. Technol. A* 14 (1996) 3267.